Synthesis and Characterization of a New Ternary Nitride, Ca₃VN₃

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We have synthesized a new ternary nitride, Ca_3VN_3 , from the binary nitrides at high temperature. The refined structure was solved in the $P2_1/m$ space group with lattice constants a = 6.717(2), b = 5.064(2), c = 6.720(3) Å, $\beta = 78.88(3)^\circ$, Z = 2, R = 3.2%, and $R_w = 3.7\%$. The structure is related to the recently reported A_3MN_3 compounds (D. A. Vennos, M. E. Badding, and F. J. DiSalvo, *Inorg. Chem.* **29**, 4059, 1990) with sheets of $[VN_3]^{6-}$ trigonal planar units and calcium ions. Ca_3VN_3 is insulating with V³⁺ in the low-spin state, S = 0. © 1992 Academic Press, Inc.

Introduction

Recently we have reported the synthesis and properties of alkaline earth-transition metal nitrides (1-3). A new structure type was observed in Ca_3CrN_3 (1), which consists of sheets of trigonal planar $[CrN_2]^{6-}$ units separated by Ca ions. This was the first solid state example of a transition metal in a trigonal planar environment of nitrogen atoms. The low Cr site symmetry (C_{2n}) and high covalency of the Cr-N bonding leads to a low-spin Cr^{3+} ($S = \frac{1}{2}$) state. Ca₃CrN₃ is the only Cr^{3+} compound known to be low spin. Since our report of Ca₃CrN₃, more alkaline earth-transition metal nitride structures have been reported to contain this $[MN_3]^{6-}$ unit; for example, Ba₃(Sr₃)FeN₃(4) and $Ca_6Fe(Ga)N_5$ (5). Although none of these structures are isotypic, they share a common structural feature: planes of $[MN_3]^{6-}$ anions separated by alkaline earth cations. However, magnetic properties of these compounds were not reported.

We report the synthesis, structure, and properties of a new compound, Ca_3VN_3 , containing $[VN_3]^{6-}$ anions which are separated by Ca^{2+} ions.

Experimental

Starting materials. Vanadium nitride was prepared by heating vanadium metal (2N5, -325 mesh) in flowing nitrogen (prepurified by passing through finely divided copper at 150°C and then activated alumina) at 1050°C for 15 hr. Granules of Ca (2N) were reacted at 1000°C for 2 days in flowing nitrogen (prepurified as above) to yield pure Ca₃N₂. The binary products were identified by X-ray powder diffraction methods. Since both Ca₃N₂ and Ca₃VN₃ are air-sensitive, all manipulations were carried out in an argonfilled glove box.

Synthesis of Ca_3VN_3 . Ca_3VN_3 was prepared as small single crystals by heating Ca granules in a vanadium foil boat to 950°C in 12 hr under Ar, soaking at 950°C for 18.5 hr under a backpressure of N₂ which allowed slow diffusion of N₂ into the Ar-filled reac-

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tion zone, followed by cooling at 22.5°/hr to 500°C, and finally cooling to room temperature in 5 hr. In order to obtain single crystals of the product, it was imperative that the heat up be under Ar and that the reaction take place slowly by indiffusion of N₂. Crystals were not obtained unless we strictly adhered to this heating schedule. The reaction yielded a distribution of brown needle crystals up to 0.5 mm in length which were used in the structure determination. Microprobe analysis of the crystals indicated the presence of Ca and V at a Ca : V ratio of approximately 3.5 : 1.

The compound was then prepared as a polycrystalline phase by heating a pressed pellet of an intimate mixture of Ca₃N₂ (0.391 g, 2.64 mmol) and VN (0.172 g, 2.64 mmol) in a sealed molybdenum tube at 1350°C for 30 hr. X-ray powder diffraction indicated a nearly pure Ca₃VN₃ with a small amount of unidentified impurity phase. The intensity of the strongest impurity diffraction peak in our best material was 13% as intense as the strongest peak of the majority phase. Other preparative conditions always produced larger quantities of the impurity phase. This product was used for both magnetic susceptibility and electrical resistivity measurements.

Structure determination. Unit cell symmetry and approximate lattice constants were obtained from axial photographs of a brown crystal mounted along the needle axis. A $0.25 \times 0.045 \times 0.018$ -mm³ crystal sealed in a 0.3-mm glass capillary under argon was used in the structure determination. Data were collected on Syntex $P2_1$ 4 circle diffractometer using Mo $K\alpha$ radiation (0.71069 Å) and a graphite monochromator. The unit cell dimensions were refined using 25 independent reflections in the 2θ range 15°-25°. An appreciable decay of approximately 25% was observed during data collection according to three check reflections measured every 50 reflections. This decay was monotonic in time and was likely due to the slow decomposition of Ca₃VN₃ in the X-ray beam. The data were scaled appropriately to correct for crystal decomposition using XDISK (6). The choice of space groups was reduced to $P2_1$ and $P2_1/m$ by systematic absences, and the structure was solved in the $P2_1/m$ space group. Refinement in the noncentrosymmetric space group resulted in no significant change in the structure or *R* factors. An empirical absorption correction was applied (the psi scan) and after averaging symmetry related reflections, 468 unique reflections with $F > 3\sigma(F)$ were used to solve the structure.

The structure determination and refinement were performed using Nicolet SHELXTL Plus (6) software running on a Microvax computer. The function minimized in the least squares refinement was $\sum w(|F_o| - |F_c|)^2$ with $w = \sigma^{-2}$. Anisotropic refinement of all atoms (43 parameters) converged to R = 3.2% and $R_w = 3.7\%$. Table I summarizes data collection parameters. Atomic positions are listed in Table II.

After the structure was solved, the powder patterns taken on a Scintag XDS2000 diffractometer could be indexed. The observed pattern minus the impurity peaks matched that calculated by Lazy Pulverix (7) on the basis of the single crystal data.

Electrical properties. The resistance of a pellet of Ca₃VN₃ was measured in a small press inside the glove box. The two pistons were electrically insulated from each other, allowing a two-point measurement of the resistance. The resistance determined on a pellet 4.5 mm in diameter and approximately 3 mm thick was larger than 30 MΩ (the upper limit of our measurement apparatus) yielding a lower limit to the resistivity on the order of $10^7 \Omega$ cm.

Magnetic susceptibility. The magnetic susceptibility of the sample was measured by the Faraday technique as previously described (8). The susceptibility of a polycrystalline sample sealed in a thin-walled high-purity quartz tube was determined to be field

TABLE I SUMMARY OF CRYSTAL AND DIFFRACTION

DATA FOR Ca₃VN₃

Space group	$P2_{1}/m$
z	2
a, b, c, (Å)	6.717(2), 5.064(2), 6.720(3)
β (deg)	78.88(3)
$V(Å^3)$	224.3(2)
Density, calc (g/cm ³)	3.157
T (K) data collection	298
Crystal dimensions (mm)	0.25 imes 0.045 imes 0.018
Absorption coeff (mm ⁻¹)	5.334
$2\theta \max (deg)$, scan type	55, $\omega - 2\theta$
Octants collected	hkl; -h-kl; -hkl; h-kl
X-ray radiation	ΜοΚα
Monochromator	Graphite
Measured reflections	1109
Observed reflections ^a	468
Independent reflections	572
F (000)	208
Number of parameters	43
Largest diff peak (e ⁻ /Å ³)	0.73
R^{b}, R_{w}^{c} (%)	3.2, 3.7

^a $F_{o} > 3\sigma(F_{o})$.

$${}^{b}R = \Sigma(|F_{o}| - |F_{c}|)/\Sigma(|F_{o}|).$$

 ${}^{c}R_{w} = [\Sigma(w(|F_{o}| - |F_{c}|)^{2})/\Sigma(w|F_{o}|^{2})]^{1/2}, w = \sigma(F_{o})^{-2}.$

independent at room temperature, indicating that no ferromagnetic impurities were present. Figure 1 shows the results of a temperature-dependent study between 4 and 300 K at a magnetic field strength of 13.6 kG.

TABLE II Positional Parameters for Ca₃VN₃

Atom	Site	x	у	z
v		0.8028(2)	0.25	0.3028(2)
Ca(A)	2e	0.3915(2)	0.25	0.8915(2)
Ca(B)	2e	0.9021(2)	0.25	0.8297(2)
Ca(C)	2 <i>e</i>	0.3298(2)	0.25	0.4022(2)
N(A)	2 <i>e</i>	0.6268(9)	0.25	0.1271(9)
N(B)	2 <i>e</i>	0.3180(10)	0.75	0.4315(9)
N(C)	2e	0.9317(10)	0.75	0.8159(9)

Results and Discussion

The structure of Ca_3VN_3 (Fig. 2) consists of trigonal planar $[VN_3]^{6-}$ anions separated by calcium ions. A view down the unique axis clearly shows the stacking arrangement of the eclipsed units. Neighboring $[VN_3]^{6-}$ anions are displaced by $\frac{1}{2}$ in z and the orientation of the triangle is reversed (Fig. 3). Although the three V–N bond distances are within 0.016 Å of each other, the bond angles reduce the symmetry of the anion to $C_{2\nu}$, similar to Ca_3CrN_3 . A list of relevant bond distances and angles can be found in Table III.



FIG. 1. Temperature dependence of the magnetic susceptibility of Ca₃VN₃.



FIG. 2. A view of Ca_3VN_3 down the unique axis. The Ca atoms are dotted circles, the V atoms are open circles, and N atoms are completely hatched circles.

The average V–N distance (1.83 Å) is close to the Cr–N distance in Ca₃CrN₃, indicating multiple metal–nitrogen bonding. Each Ca atom is coordinated by three nitrogen atoms and two calcium atoms in the same plane, and by two nitrogen atoms above and below that plane at 2.54 \pm 0.01 Å. The three coordinated nitrogen atoms subtend an angle of 170°C, creating a quite asymmetric environment about the central



FIG. 3. A view perpendicular to the unique axis shows the stacking arrangement and the labels used to identify atoms in Table III. Some atoms in the unit cell are not included for clarity.

Ca. The in-plane Ca–N distances for each of the three inequivalent Ca's are: Ca(A) 2.443, 2.625, and 2.634 Å; Ca(B) 2.447, 2.503, and 2.818 Å; Ca(C) 2.445, 2.494, and 2.804 Å. The shorter Ca–N distances are comparable to those observed in Ca₃N₂ (2.46 Å), Ca₃BiN (2.45 Å) (3), and CaNiN (2.50 Å) (2). The closest Ca–Ca distance between planar units (3.39 Å) is slightly longer than that in Ca₃N₂ (3.16 Å) (9).

The V-V distances are much too long for

TABLE III

Bond distances		Bond angles		
V(1) - V(2)	5.086(2)	N(1C)-V(1)-N(1A)	114.8(3)	
V(1) - V(3)	5.064(2)	N(1C) - V(1) - N(1B)	130.7(3)	
V(1)-Ca(2C)	3.232(2)	N(1B) - V(1) - N(1B)	114.8(3)	
V(1)-Ca(2A)	3.237(2)		. ,	
V(1)-Ca(2B)	3.233(2)			
V(1) - N(1B)	1.820(6)			
V(1)–N(1A)	1.827(7)			
V(1)–N(1C)	1.811(6)			
Ca(2A)-N(2B)	2.616(6)			
Ca(2A) - N(1A)	2.539(1)			
Ca(2A) - N(1A)	2.443(6)			
Ca(1C)-N(1A)	2.445(6)			

IMPORTANT DISTANCES (Å) AND ANGLES (Degrees) IN Ca₃VN₃

any metal-metal interactions (shortest distance 5.06 Å from V1 to V3 in Fig. 3). As expected from the long V-V distances, Ca_3VN_3 is insulating. The susceptibility can be fit with a simple Curie-Weiss expression,

$$\chi_{\rm g} = \chi_{\rm o} + C_{\rm g}/T + \theta,$$

and we obtain $\chi_0 = 1.24 \times 10^{-7}$ emu/g, $C_{\rm g} = 9.04 \times 10^{-6} \, {\rm emu/g}$, and $\theta = 5 \, {\rm K}$ (with a mean square deviation of 2.1%). The small value of C_{g} indicates that the temperature dependence arises from a paramagnetic impurity. If the V had a moment appropriate to V^{3+} (S = 1), the value of C_g would be 4.70×10^{-3} emu/g. We conclude that V in Ca_3VN_3 and in the impurity phase present in the sample must be nonmagnetic. Since an unknown impurity is present, a definitive value of χ_0 for Ca₃VN₃ cannot be determined from this measurement. The S = 0assignment for V^{3+} in Ca₃VN₃ is consistent with the low site symmetry and the low-spin configuration in Ca₃CrN₃.

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